1. (15 points) If \hat{A}_s is an operator in the Shrödinger picture and \hat{A}_H is the corresponding operator in the Heisenberg picture, Such that

$$\hat{A_H} = e^{\frac{-i\hat{H}t}{\hbar}} \hat{A_s} e^{\frac{i\hat{H}t}{\hbar}}$$

Under what conditions $\hat{A_H}$ can be time-independent. Justify your answer. A_s must be time independent and commute with the Hamiltonian.

2. A Hamiltonian H has two orthonormal eigenstates |1 > and |2 > such that:

$$\hat{H}|1>=E_1|1>$$
 $\hat{H}|2>=E_2|2>$ $E_1\neq E_2$

Two states |A > and |B > are defined as follows:

$$<1|A>=\frac{1}{\sqrt{2}} \qquad <2|A>=\frac{i}{\sqrt{2}} \\ <1|B>=\frac{1}{\sqrt{2}} \qquad <2|B>=\frac{-i}{\sqrt{2}}$$

(a) (6 points) Calculate $\langle A|B \rangle$ and $\langle B|A \rangle$

$$\begin{array}{rcl} < A|B> &=& < A|I|B> = < A|(|1><1|+|2><2|)|B> \\ &=& < A|1><1|B>+<2|B> \\ &=& \frac{1}{\sqrt{2}}\times\frac{1}{\sqrt{2}}+\frac{-i}{\sqrt{2}}\times\frac{-i}{\sqrt{2}}=0 \end{array}$$

(b) (6 points) If the system initially in the state $|\psi(t=0)\rangle = |A\rangle$, what is the time dependent state $|\psi(t)\rangle$?

$$\begin{split} |\psi(t=0)>&=|A>=|1>+|2>\\ |\psi(t)>&=|1>e^{-iE_{1}t/\hbar}+|2>e^{-iE_{2}t/\hbar} \end{split}$$

(c) (8 points) What is the probability of finding the particle at state $|A\rangle$ at any time t? Probability is defined as:

 $|\langle A|\psi(t)\rangle|^2 = |e^{-iE_1t/\hbar} + e^{-iE_2t/\hbar}|^2$

3. Consider an electron in the Hydrogen-atom. The wavefunction of the electron is, at time t = 0, written as:

$$\Psi(r,t=0) = A(\psi_{211} + 2\psi_{300} + \psi_{421})$$

(a) (4 points) Find the normalization constant A Since ψ_{nlm} forms an orthonormal set then:

$$\begin{array}{rcl}
1 & = & A^2(1+4+1) \\
A & = & \frac{1}{\sqrt{6}}
\end{array}$$

(b) (4 points) Write the wavefunction at any later time t

$$\Psi(r,t=0) = \frac{1}{\sqrt{6}} (\psi_{211}e^{-iE_2t/\hbar} + 2\psi_{300}e^{-iE_3t/\hbar} + \psi_{421}e^{-iE_4t/\hbar})$$

(c) (4 points) What is the expectation value of L_z . Since ψ_{nlm} is an eigenvector of L_z , then

$$< L_z > = \sum_{n=1}^{\infty} c_n^2 L_z$$

= $\frac{1}{6}(\hbar) + \frac{4}{6}(0) + \frac{1}{6}(\hbar) = \hbar/3$

(d) (4 points) What is the expectation value of L^2 Since ψ_{nlm} is an eigenvector of L^2 , then

$$< L^2 > = \sum_{l=0}^{\infty} c_n^2 \hbar^2 l(l+1)$$

= $\frac{1}{6} (2\hbar^2) + \frac{4}{6} (0\hbar^2) + \frac{1}{6} (6\hbar^2) = 4\hbar^2/3$

(e) (4 points) What is the expectation value of H. Since ψ_{nlm} is an eigenvector of H, then

$$\langle H \rangle = \sum_{n=1}^{\infty} c_n^2 E_n$$

= $\frac{1}{6}(E_2) + \frac{4}{6}(E_3) + \frac{1}{6}(E_4)$

- 4. Consider a particle with spin-angular momentum j=3/2. There are four sub-levels with this value of j but different eigenvalues of \hat{j}_z . They are written in the following form $|jm_j\rangle$: $|\frac{3}{2}\frac{3}{2}\rangle$, $|\frac{3}{2}\frac{1}{2}\rangle$, $|\frac{3}{2}\frac{-1}{2}\rangle$, and $|\frac{3}{2}\frac{-3}{2}\rangle$.
 - (a) (6 points) Show that the raising operator can be written in the following form:

$$\hat{j}_{+} = \sqrt{3} |\frac{3}{2} \frac{3}{2} > < \frac{3}{2} \frac{1}{2}| + 2|\frac{3}{2} \frac{1}{2} > < \frac{3}{2} \frac{-1}{2}| + \sqrt{3} |\frac{3}{2} \frac{-1}{2} > < \frac{3}{2} \frac{-3}{2}|$$

We know(can calculate) the following relations:

$$\begin{aligned} j_{+} |\frac{3}{2}\frac{3}{2} > &= 0 \\ j_{+} |\frac{3}{2}\frac{1}{2} > &= \sqrt{3}|\frac{3}{2}\frac{3}{2} > \\ j_{+} |\frac{3}{2}\frac{-1}{2} > &= 2|\frac{3}{2}\frac{1}{2} > \\ j_{+} |\frac{3}{2}\frac{-3}{2} > &= \sqrt{3}|\frac{3}{2}\frac{-1}{2} > \\ j_{+} |\frac{3}{2}\frac{-3}{2} > &= \sqrt{3}|\frac{3}{2}\frac{-1}{2} > \\ j_{+} &= j_{+}(|\frac{3}{2}\frac{3}{2} > < \frac{3}{2}\frac{3}{2}| + |\frac{3}{2}\frac{1}{2} > < \frac{3}{2}\frac{1}{2}| + |\frac{3}{2}\frac{-1}{2} > < \frac{3}{2}\frac{-1}{2}| + |\frac{3}{2}\frac{-3}{2} > < \frac{3}{2}\frac{-3}{2}|) \end{aligned}$$

and now we can get the required results.

(b) (3 points) Write \hat{j}_{-} We know that $j_{-} = j_{+}^{\dagger}$

$$\hat{j}_{+} = \sqrt{3}|\frac{3}{2}\frac{1}{2}> <\frac{3}{2}\frac{3}{2}|+2|\frac{3}{2}\frac{-1}{2}> <\frac{3}{2}\frac{1}{2}|+\sqrt{3}|\frac{3}{2}\frac{-3}{2}> <\frac{3}{2}\frac{-1}{2}|$$

(c) (8 points) Find the matrix representation of \hat{j}_x , \hat{j}_y , and \hat{j}_z

$$\hat{j}_{z} = \hbar \begin{pmatrix} 3/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & -1/2 & 0 \\ 0 & 0 & 0 & -3/2 \end{pmatrix}$$
$$\hat{j}_{+} = \hbar \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & \sqrt{3} \\ 0 & 0 & 0 & 0 \end{pmatrix}$$
$$\hat{j}_{-} = \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ \sqrt{3} & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}$$
$$\hat{j}_{x} = \frac{\hat{j}_{+} + \hat{j}_{-}}{2} = \frac{\hbar}{2} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ \sqrt{3} & 0 & 2 & 0 \\ 0 & 2 & 0 & \sqrt{3} \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}$$
$$\hat{j}_{x} = \frac{\hat{j}_{+} + \hat{j}_{-}}{2} = \frac{\hbar}{2} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ \sqrt{3} & 0 & 2 & 0 \\ 0 & 2 & 0 & \sqrt{3} \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}$$

(d) (4 points) let $|\psi(t=0)\rangle = \frac{1}{2\sqrt{2}}(|\frac{3}{2}\frac{3}{2}\rangle + \sqrt{3}|\frac{3}{2}\frac{1}{2}\rangle + \sqrt{3}|\frac{3}{2}\frac{1}{2}\rangle + |\frac{3}{2}\frac{-3}{2}\rangle)$ Show that it is an eigenvector of j_x . One can write the wavefunction in matrix form as:

$$\hat{j}_x == \frac{1}{2\sqrt{2}} \begin{pmatrix} 1\\\sqrt{3}\\\sqrt{3}\\1 \end{pmatrix}$$

Now one can multiply the matrix of \hat{j}_x with the wave function and show that this wave function is indeed an eignevector \hat{j}_x with an eigenvalue of 3/2

- (e) (4 points) If the particle a magnetic moment $\vec{\mu} = g\vec{j}$, and is placed in uniform magnetic field that points in the x-direction. The particle initial wave-function is given in the previous part. Find $\langle \hat{j}_z \rangle(t)$ One can notice that the Hamiltonian is probational to j_x and thus the wavefunction in the previous part is a stationary state. Thus the expectation value of any operator doesn't depend on time. $\langle j_z \rangle(t) = 0$
- 5. (20 points) Show that any operator that commutes with two Cartesian components of the angular momentum operator necessarily commutes with the total angular momentum operator. We are given the following:

$$[A, L_i] = [A, L_j] = 0$$

All what we need now is to show that

$$[A, L_k] = 0$$

We know that

$$[L_i, L_j] = i\hbar L_k$$

Now:

$$[A, L_z] = \left[A, \frac{i}{\hbar} [L_i, L_j]\right]$$
$$= \frac{i}{\hbar} [A, L_i L_j - L_j L_i]$$
$$= \frac{i}{\hbar} ([A, L_i L_j] - [A, L_j L_i])$$
$$= 0$$

Question:	1	2	3	4	5	Total
Points:	15	20	20	25	20	100
Score:						
Good Luck						